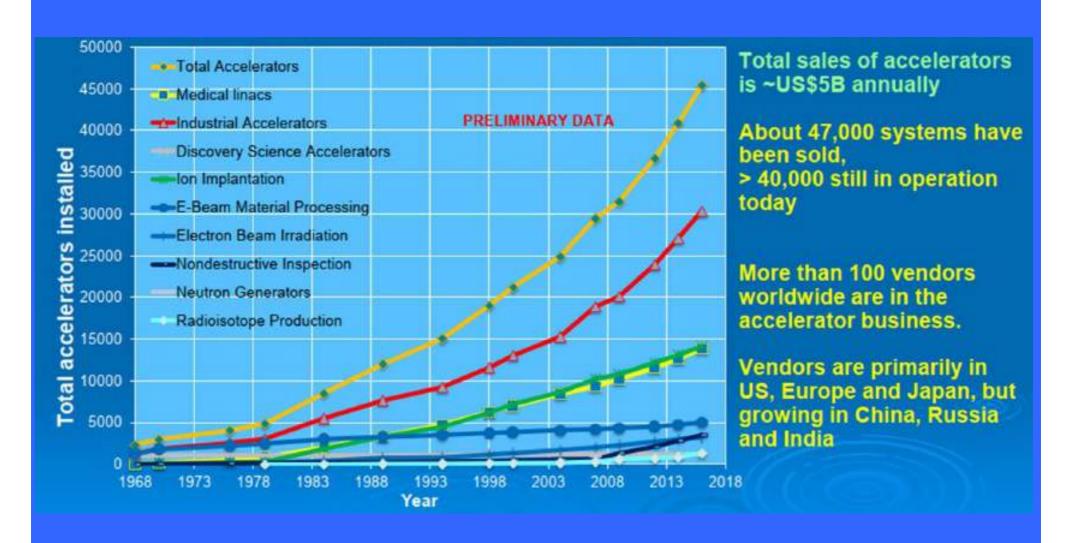
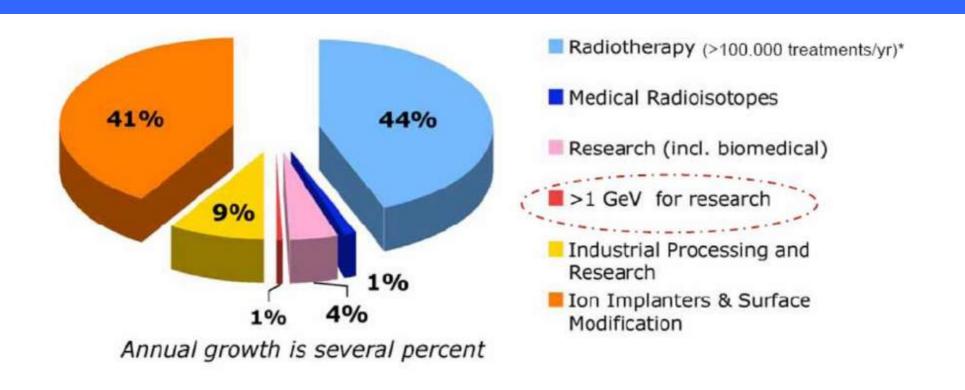
# Introduzione alla Fisica degli Acceleratori di Particelle Massimo.Ferrario@LNF.INFN.IT

### Accelerators installed worldwide

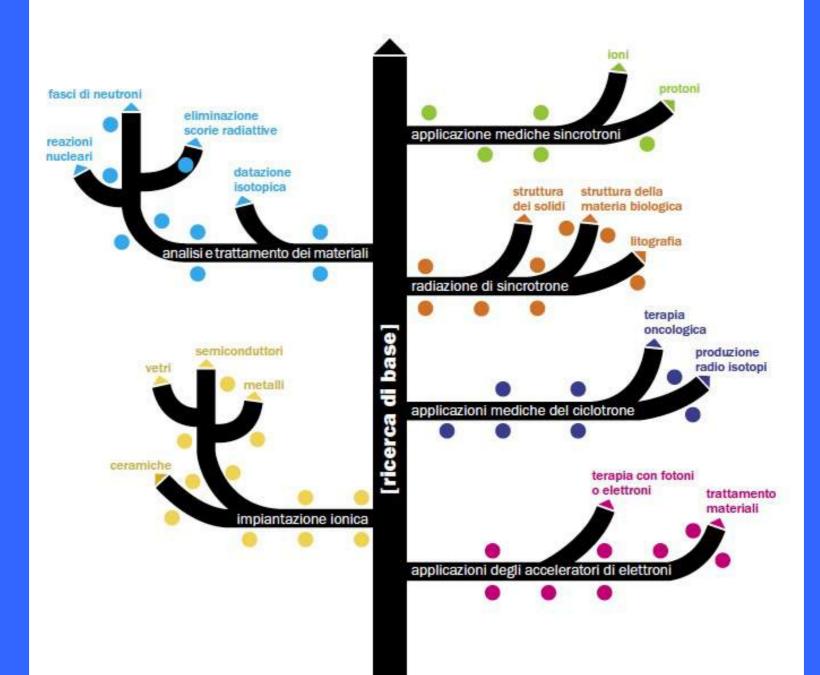


-Accelerators for Americas Future Report, pp. 4, DoE, USA, 2011

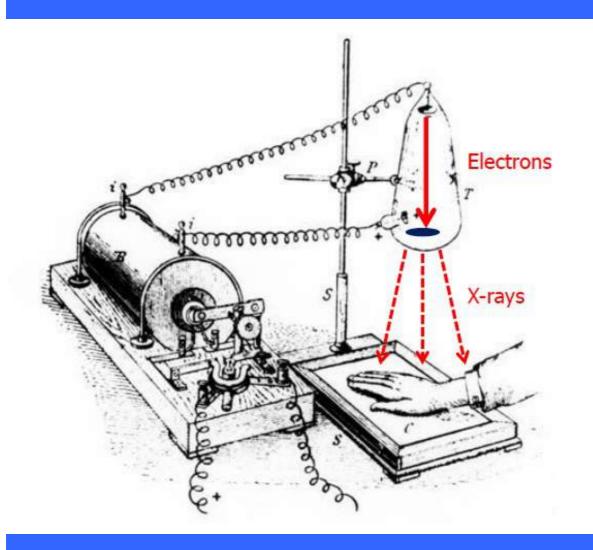
### Accelerators installed worldwide

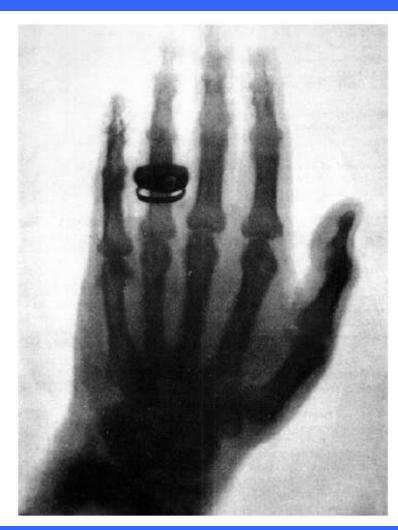


-Accelerators for Americas Future Report, pp. 4, DoE, USA, 2011

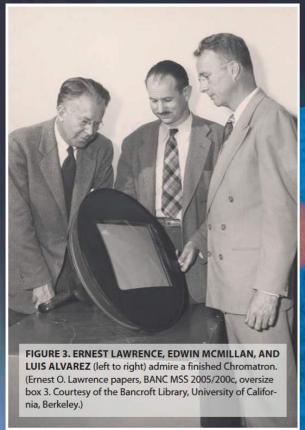


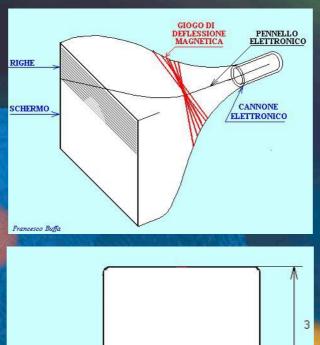
# Roentgen 1896 – First radiograph of a hand



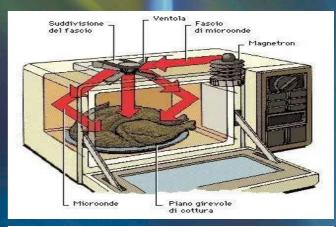


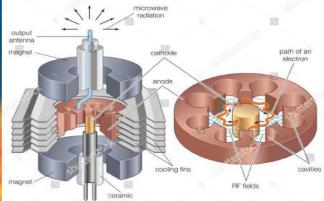
### Acceleratori Domestici





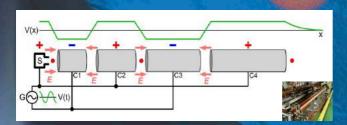
Francesco Buffa

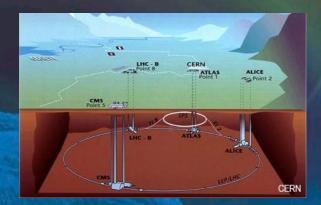


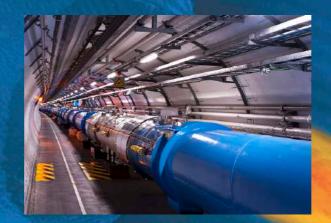


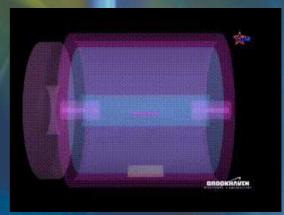
### Acceleratori ad Alta Energia









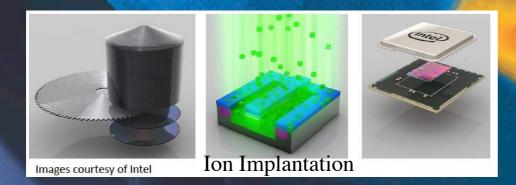


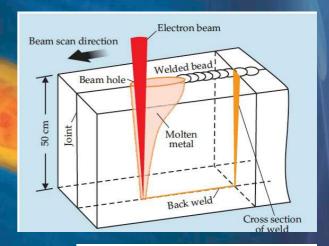


### Acceleratori per l'Industria





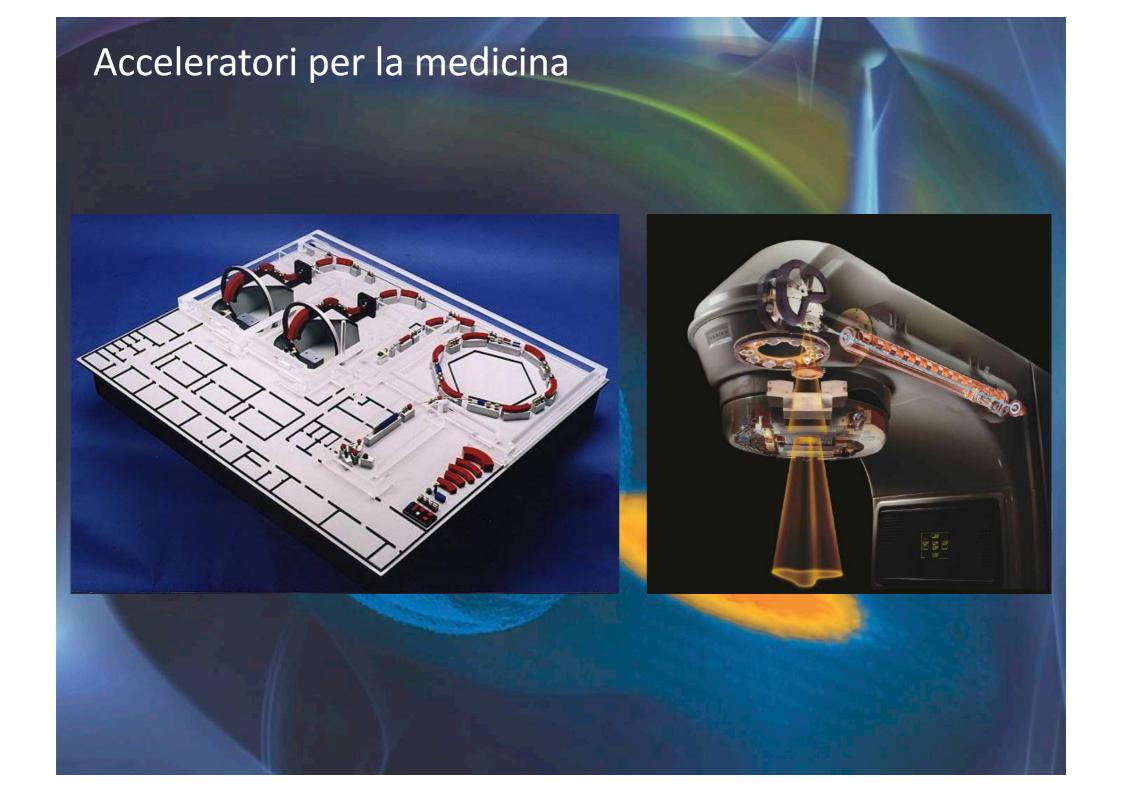




The beam business: Accelerators in industry

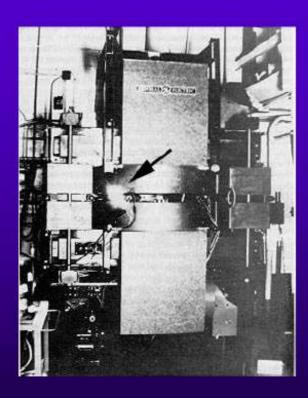
Robert W. Hamm, and Marianne E. Hamm

Citation: Physics Today 64, 6, 46 (2011); doi: 10.1063/1.3603918



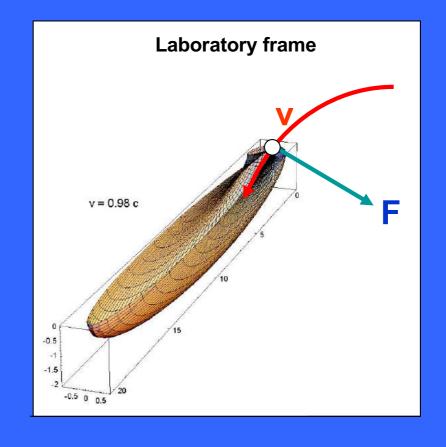
# **Synchrotron Radiation**

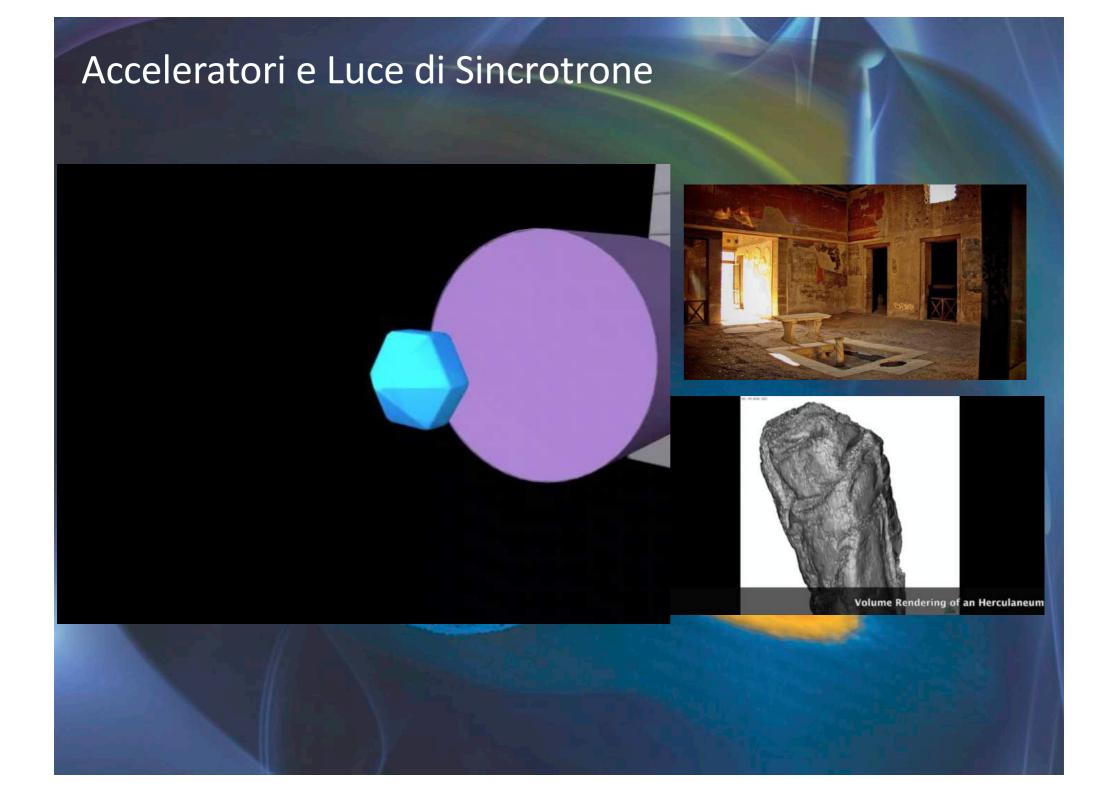
# **GE Synchrotron New York State**

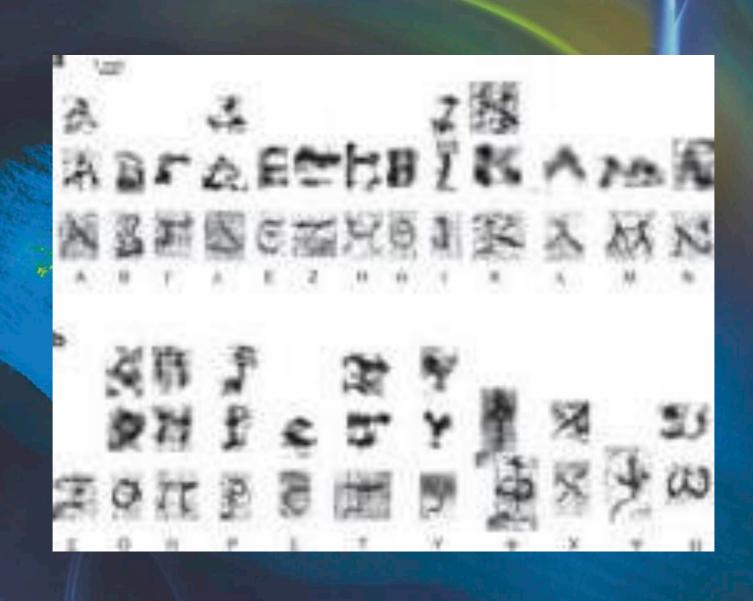


First light observed 1947

$$P_{\gamma} = \frac{cC_{\gamma}}{2\pi} \cdot \frac{E^4}{\rho^2}$$





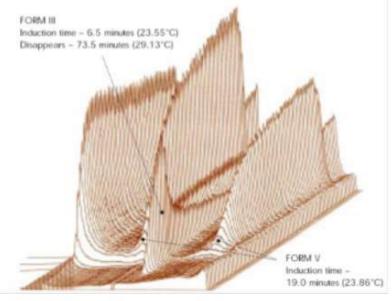


# Prevent Chocolate Melting

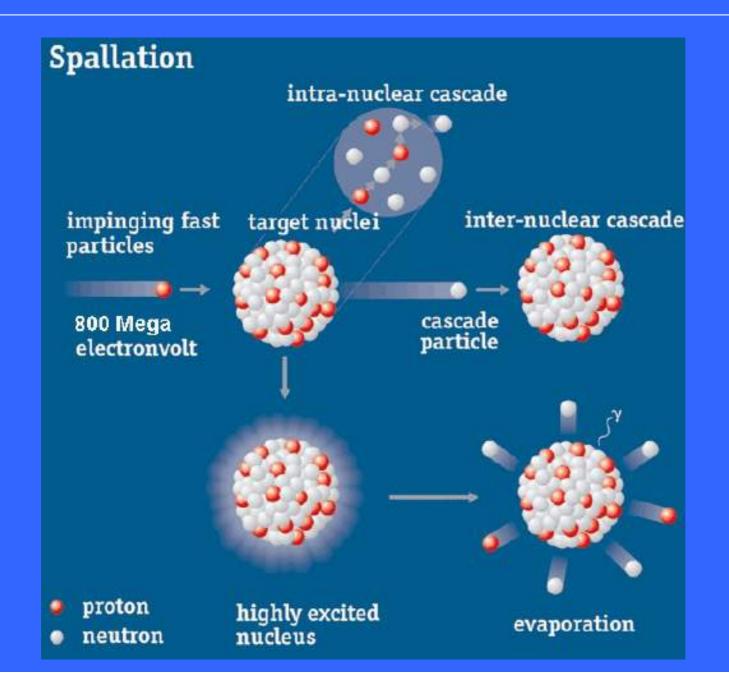


Of the six possible crystal forms, the fifth (form V) produces the best quality chocolate

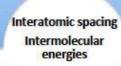
Cadbury used X-rays from a particle accelerator to study how cocoa crystallises



### **Neutron Sources**



### **Relevant Neutron Properties**

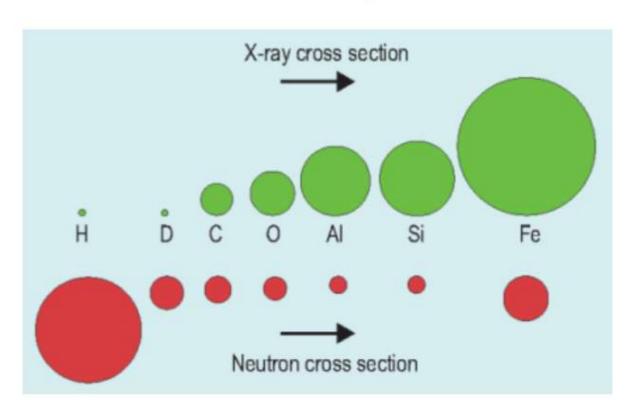


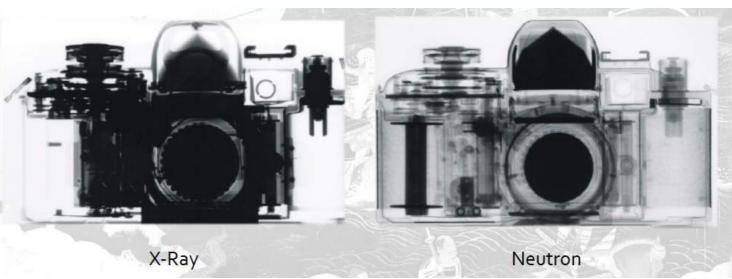
Highly penetrating (neutral particle)

> Isotopic substitution (high contrast possible e.g. H/D)

Scattering by magnetic moments

A probe of fundamental properties





The special feature of Neutrograph is it's intensity together with a moderate collimation.

These properties allow the investigation of dynamic processes with an excellent time resolution and the transmittance through strongly absorbing and bulky materials.

A totally new spectrum of scientific and engineering applications could be developed.

Among the experiments are investigations of heat exchangers and combustion engines, parts from aircrafts, fossils and historical heritage.

Institut Laue-Langevin (ILL) in Grenoble

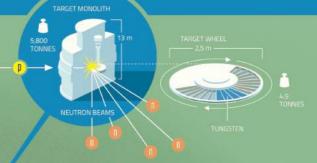


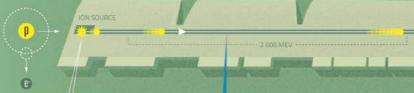
### **European Spallation Source**





occurs and neutrons are scattered from the tungsten nucleus. The more neutron produced and collected in the target, the "brighter" the neutron source. The







Electromagnetic fields are used to accelerate the protons to approximatel





Area	Application	Beam	Accelerator	Beam ener- gy/MeV	Beam current/ mA	Number
Medical	Cancer therapy	е	linac	4-20	102	>14000
		р	cyclotron, synchrotron	250	10-6	60
		С	synchrotron	4800	10-7	10
	Radioisotope production	р	cyclotron	8-100	1	1600
Industrial	lon implantation	B, As, P	electrostatic	< 1	2	>11000
	Ion beam analysis	р, Не	electrostatic	<5	10-4	300
	Material processing	е	electrostatic, linac, Rhodatron	≤10	150	7500
	Sterilisation	е	electrostatic, linac, Rhodatron	≤10	10	3000
Security	X-ray screening of cargo	е	linac	4-10	?	100?
	Hydrodynamic testing	е	linear induction	10-20	1000	5
Synchrotron light sources	Biology, medicine, materials science	е	synchrotron, linac	500-10000		70
Neutron scattering	Materials science	р	cyclotron, synchrotron, linac	600-1000	2	4
Energy - fusion	Neutral ion beam heating	d	electrostatic	1	50	10
	Heavy ion inertial fusion	Pb, Cs	Induction linac	8	1000	Under development
	Materials studies	d	linac	40	125	Under development
Energy - fission	Waste burner	р	linac	600-1000	10	Under development
	Thorium fuel amplifier	р	linac	600-1000	10	Under development
Energy - bio-fuel	Bio-fuel production	е	electrostatic	5	10	Under development
Environmental	Water treatment	е	electrostatic	5	10	5
	Flue gas treatment	е	electrostatic	0.7	50	Under development

### Ultra-precise microscopy

- Probing particles are required for studies of the elementary constituents
- The associated de Broglie wavelength  $\lambda$  of a probing particle defines the minimum object size that can be resolved.

$$\lambda = \frac{h}{p} = h \times \frac{c}{E} \quad \text{with} \begin{cases} h = 4 \times 10^{-15} \, \text{eVs (Plank constant)} \\ p = \text{momentum, } E = \text{energy} \end{cases}$$

Resolving Smaller Objects Requires Higher Momentum Probe Particles

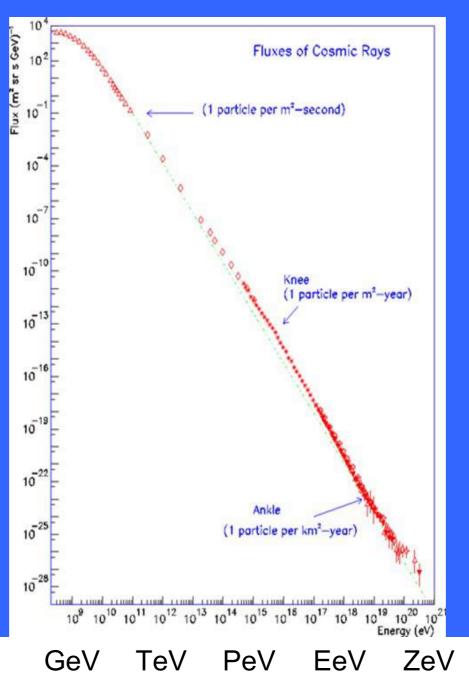
### Example of probe wavelength

- electrons with  $p = 1 \text{ keV/c} \Rightarrow /= \text{h/p} = 4 \times 10^{-12} \text{ m}$
- photons with  $E = 1 \text{ keV} \Rightarrow / = h \times c/E \sim 1.2 \times 10^{-9} \text{ m}$ .
- electrons have ~ 300 times better resolution than photons (electron-microscopy!)

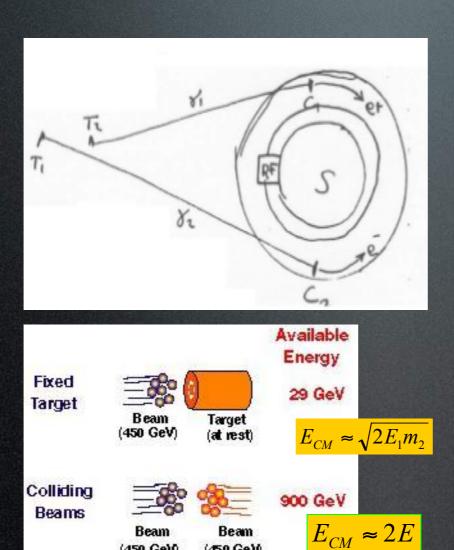
### Typical microscopic sizes

- Atom 10<sup>-10</sup> m
- Nucleus 10<sup>-14</sup> m
- Proton 10<sup>-15</sup> m
- Quark 10<sup>-19</sup> m

# Cosmos is able to accelerate particles ...but in uncontrolled conditions

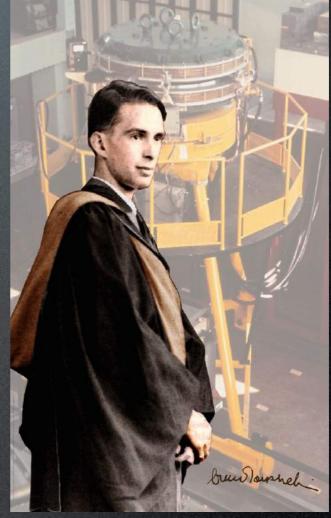


### Touschek's Anello Di Accumulazione (ADA) 1961 the first e+e- Collider



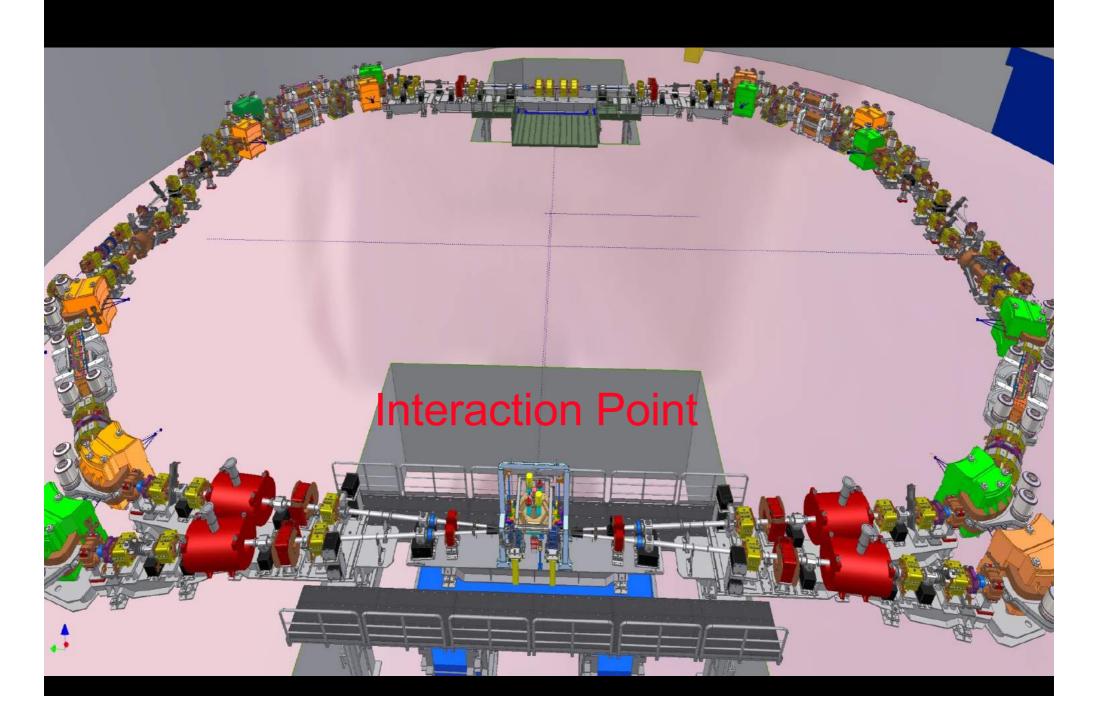
(450 GeV)

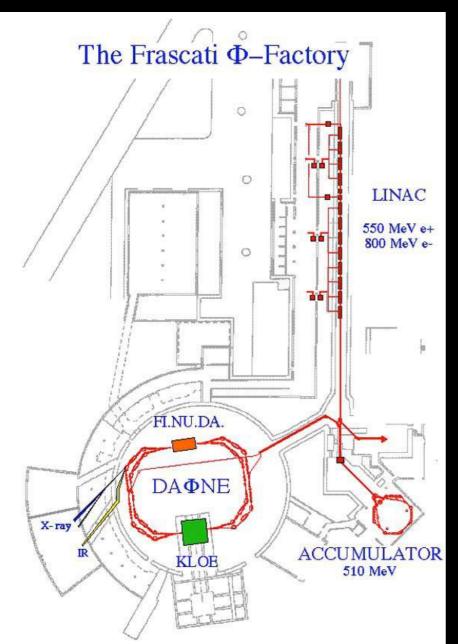
Beam (450 GeV)



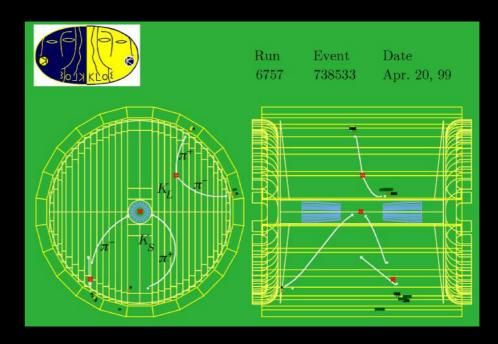


# Collider e+ e- DAFNE (INFN)









### **Historical Milestones**

- 1900 to 1925 radioactive source experiments à la Rutherford -> request for higher energy beams;
- 1928 to 1932 electrostatic acceleration ->
  - Cockcroft & Walton \* -> voltage multiplication using diodes and oscillating voltage (700 kV);

\* Nobel 1951

- Van der Graaf -> voltage charging through mechanical belt (1.2 MV);
- 1928 resonant acceleration -> Ising establish the concept, Wideroe builds the first linac;
- ♦ 1929 cyclotron ->
  - small prototype by Livingstone (PhD thesis), large scale by Lawrence\*\*;

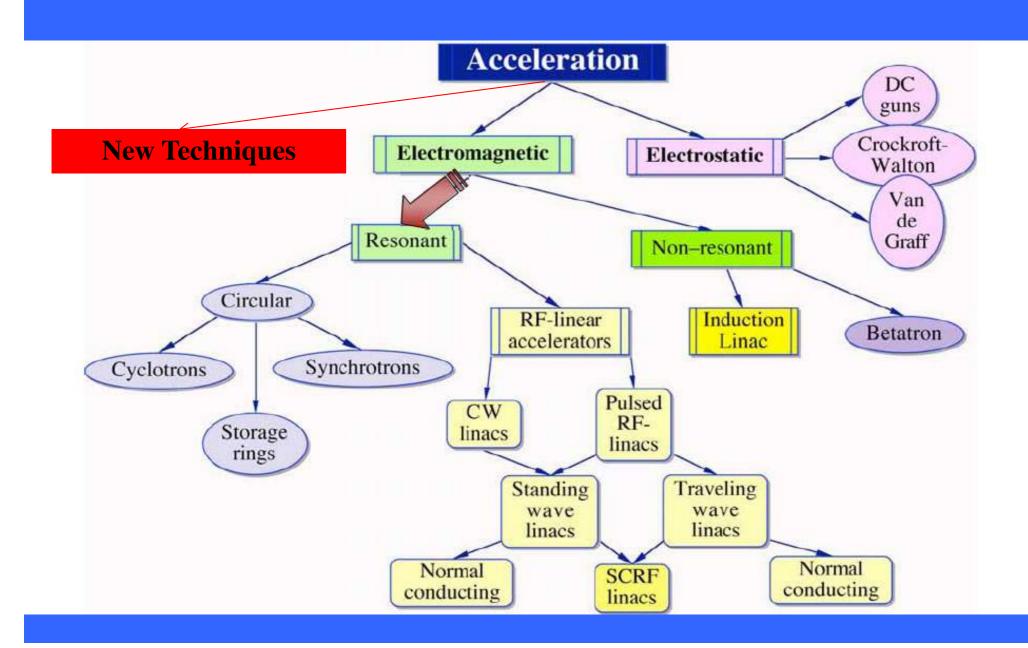
\*\* Nobel 1939

- ◆ 1942 magnetic induction -> Kerst build the betatron;
- 1944 synchrotron -> MacMillan, Oliphant and Veksel invent the RF phase stability (longitudinal focusing);
- 1946 proton linac -> Alvarez build an RF structure with drift tubes (progressive wave in  $2\pi$  mode);
- 1950 strong focusing -> Christofilos patent the alternate gradient concept (transverse strong focusing);
- 1951 tandem -> Alvarez upgrade the electrostatic acceleration concept and build a tandem;
- 1955 AGS -> Courant, Snider and Livingstone build the alternate gradient Cosmotron in Brookhaven;
- 1956 collider -> Kerst discuss the concept of colliding beams;
- ◆ 1961 e<sup>+</sup>e<sup>-</sup> collider -> Touschek invent the concept of particle-antiparticle collider;
- 1967 electron cooling -> Budker proposes the e-cooling to increase the proton beam density;
- 1968 stochastic cooling ->

\*\*\* Nobel 1984

- Van der Meer\*\*\* proposes the stochastic cooling to compress the phase space;
- 1970 RFQ -> Kapchinski & Telyakov build the radiofrequency quadrupole;
- ◆ 1980 to now superconducting magnets -> developed in various laboratories to increase the beam energy;
- ◆ 1980 to now superconducting RF -> developed in various lab to increase the RF gradient.

## **Accelerator Configurations**

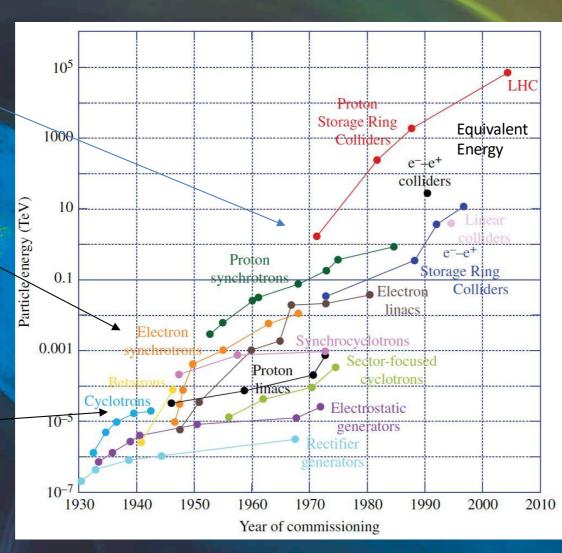


# Il diagramma di Livingstone



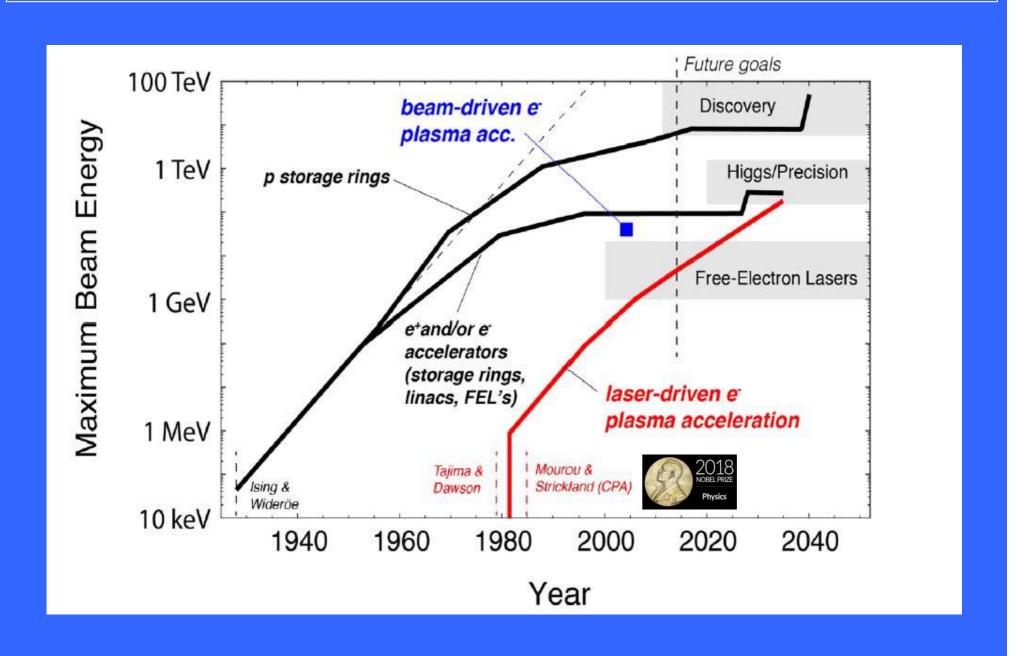




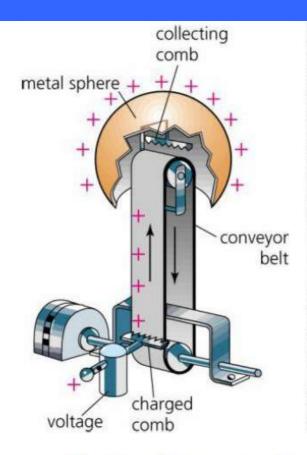


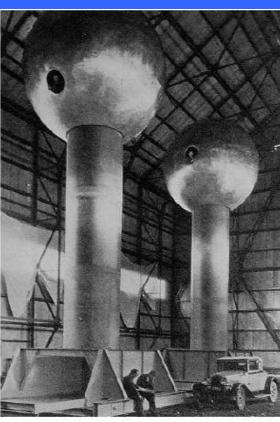
Energy of colliders is plotted in terms of the laboratory energy of particles colliding with a proton at rest to reach the same center of mass energy.

### Recent sign of saturation?



# Electrostatic Accelerator: Van de Graff





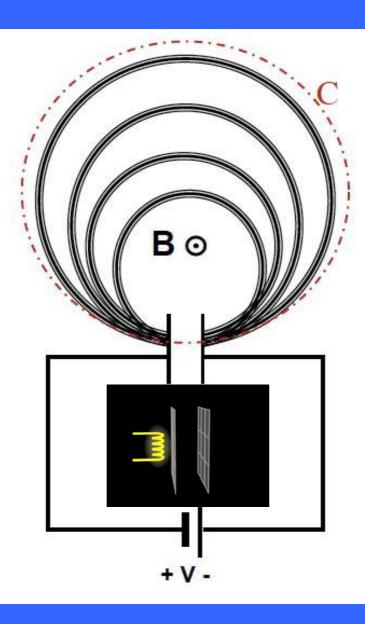


R.J. Van de Graaf

7 MV Van de Graaf at MIT (1933)

- Electric charges are transported mechanically on an insulating belt
- Stable, continuous beams, practical limit 10 15 MV

# Possible Higher energy DC accelerator?



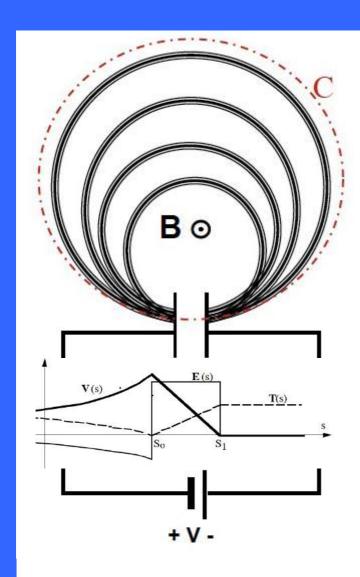
$$F_{Lorentz} = q v B = F_{centripital} = \frac{mv^2}{\rho}$$

$$\Rightarrow \rho = \frac{mv}{qB} = \frac{p}{qB}$$

$$\rho(m) = 3.34 \left(\frac{p}{1 \text{ GeV/c}}\right) \left(\frac{1}{q}\right) \left(\frac{1 \text{ T}}{B}\right)$$

$$T=q\Delta V$$

# Forbidden by Maxwell



$$\nabla \times \mathbf{E} = -\frac{d\mathbf{B}}{dt}$$

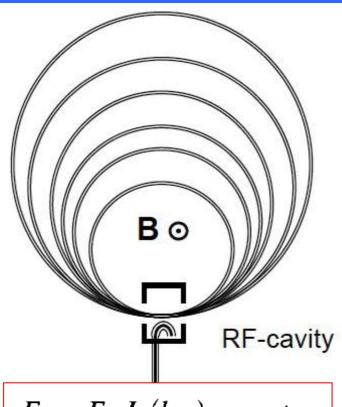
or in integral form

$$\oint_C \mathbf{E} \cdot d\mathbf{s} = -\frac{\partial}{\partial t} \int_S \mathbf{B} \cdot \mathbf{n} \, da$$

... There is no acceleration without time-varying magnetic flux

$$\Delta V_T = 0$$

# B can vary in a RF cavity



$$E_z = E_0 J_0(k_r r) \cos \omega t$$

$$B_{\theta} = -\frac{E_0}{c} J_1(k_r r) \sin \omega t$$

Note that inside the cavity  $dB/dt \neq 0$ 

However,

**Synchronism condition:** 

$$\Delta \tau_{\rm rev} = N/f_{\rm rf}$$

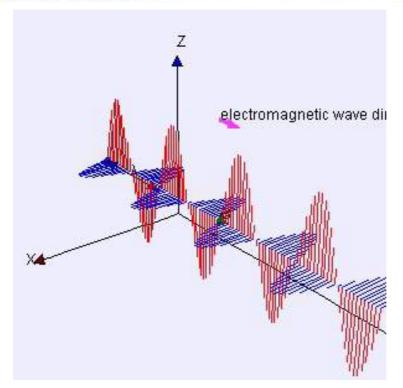
### **Lawson-Woodward Theorem**

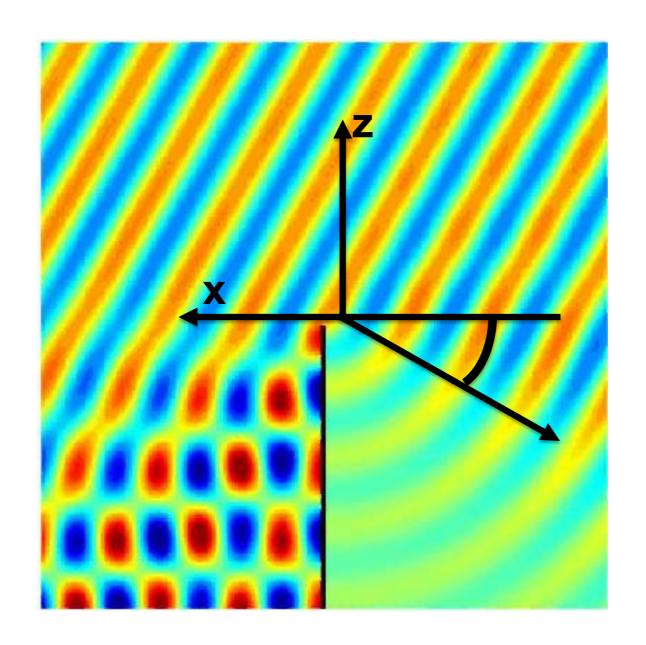
(J.D. Lawson, IEEE Trans. Nucl. Sci. NS-26, 4217, 1979)

The net energy gain of a relativistic electron interacting with an electromagnetic field in vacuum is zero.

### The theorem assumes that

- (i) the laser field is in vacuum with no walls or boundaries present,
- (ii) the electron is highly relativistic (v ≈ c) along the acceleration path,
- (iii) no static electric or magnetic fields are present,
- (iv) the region of interaction is infinite,

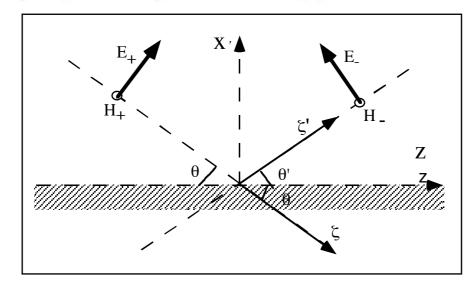




### Reflection of plane waves

Plane wave reflected by a perfectly conducting plane

$$\sigma = \infty$$



In the plane xz the field is given by the superposition of the incident and reflected wave:

$$E(x,z,t) = E_{+}(x_o,z_o,t_o)e^{i\omega t - ik\xi} + E_{-}(x_o,z_o,t_o)e^{i\omega t - ik\xi'}$$

$$\xi = z \cos \theta - x \sin \theta$$
  $\xi' = z \cos \theta' + x \sin \theta'$ 

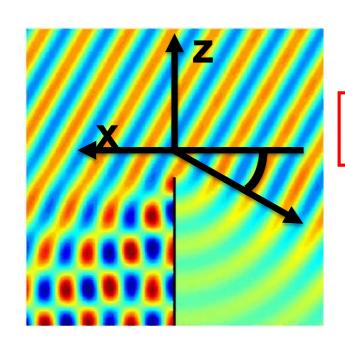
And it has to fulfill the boundary conditions (no tangential E-field)

### Reflection of plane waves (a first boundary value problem)

Taking into account the boundary conditions the longitudinal component of the field becomes:

$$E_{z}(x,z,t) = (E_{+}\sin\theta)e^{i\omega t - ik(z\cos\theta - x\sin\theta)} - (E_{+}\sin\theta)e^{i\omega t - ik(z\cos\theta + x\sin\theta)}$$

 $=2iE_{+}\sin\theta\sin(kx\sin\theta)e^{i\omega t-ikz\cos\theta}$ 



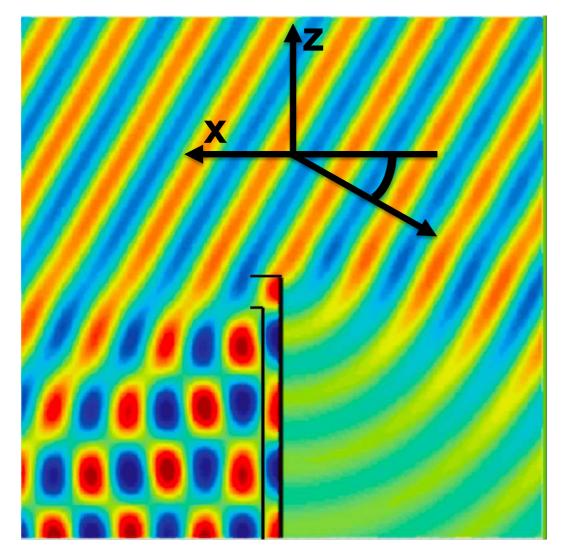
Standing Wave pattern (along x)

Guided wave pattern (along z)

The phase velocity is given by

$$v_{\phi z} = \frac{\omega}{k_z} = \frac{\omega}{k \cos \theta} = \frac{c}{\cos \theta} > c$$

## From reflections to waveguides



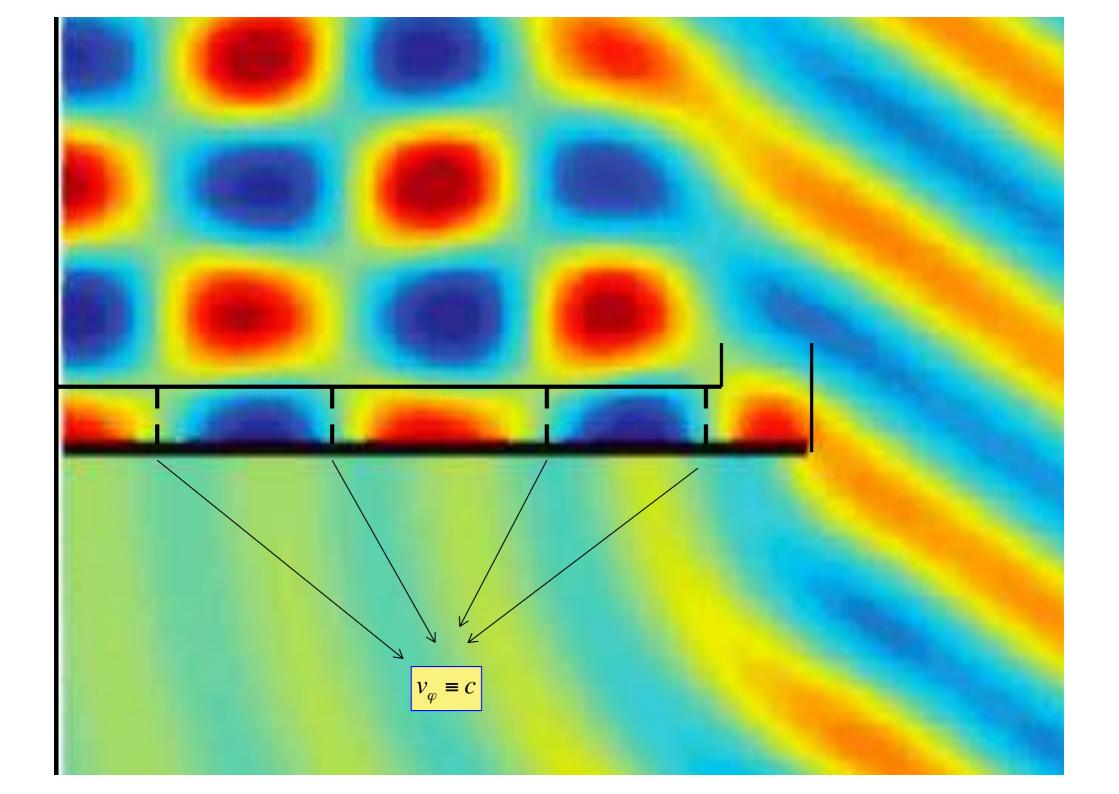
Put a metallic boundary where the field is zero at a given distance from the wall.

Between the two walls there must be an integer number of half wavelengths (at least one).

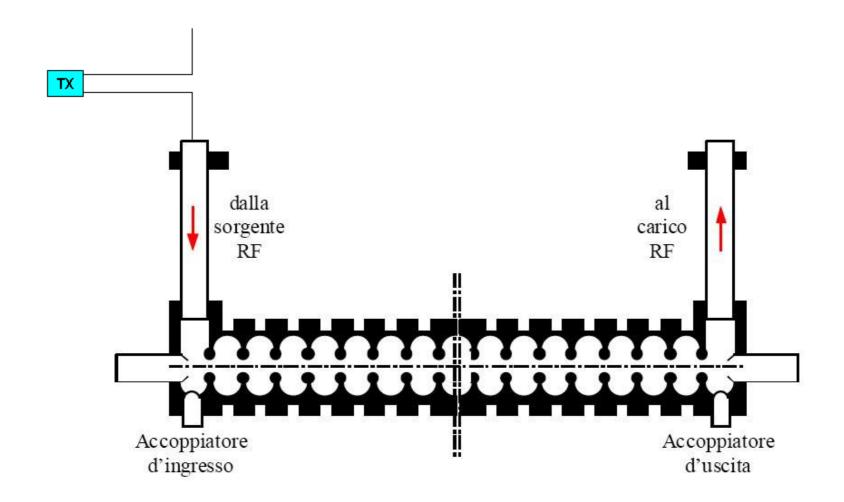
For a given distance, there is a maximum wavelength, i.e. there is **cut-off frequency**.

$$v_{\phi z} = \frac{\omega}{k} = \frac{\omega}{k \cos \theta} = \frac{c}{\cos \theta} > c$$

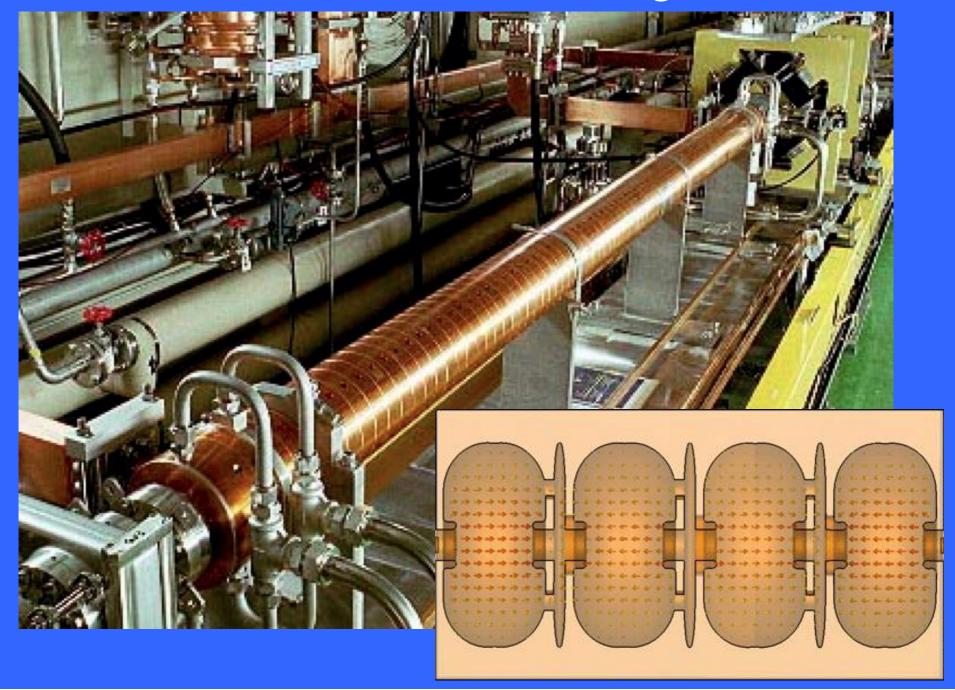
It can not be used as it is for particle acceleration



## Electromagnetic waveguides



# Conventional RF accelerating structures



We must slow down the wave propagation

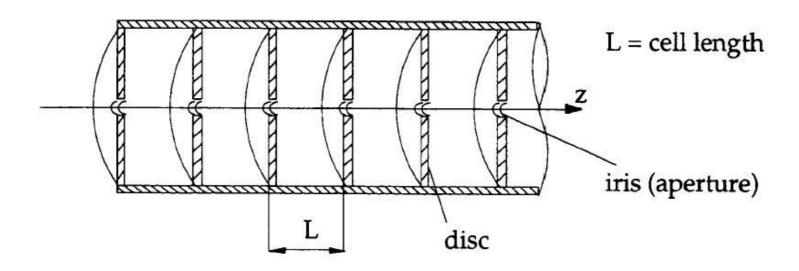
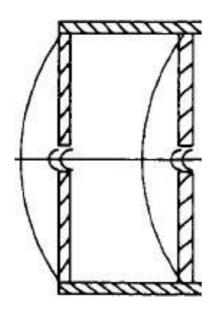


Fig. 6 Disc-loaded cavity (schematic)

In order to slow down the waves we have to load the cavity by introducing some periodic obstacle into it

$$\frac{\partial^2 E_z}{\partial z^2} + \frac{1}{r} \frac{\partial E_z}{\partial r} + \frac{\partial^2 E_z}{\partial r^2} - \frac{1}{c^2} \frac{\partial^2 E_z}{\partial t^2} = 0$$



$$E_z = E_0 J_0(k_r r) \cos \omega t$$

$$B_{\theta} = -\frac{E_0}{c} J_1(k_r r) \sin \omega t$$

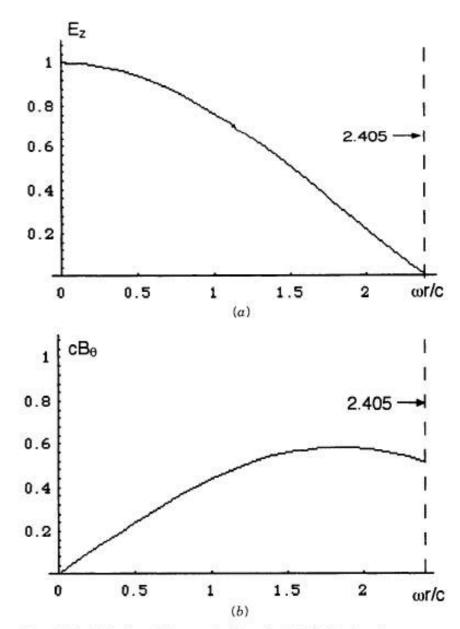
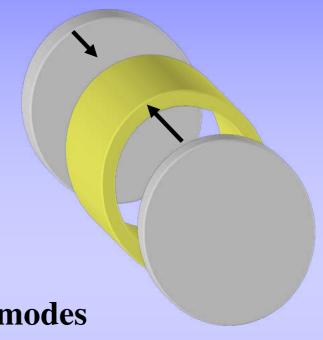


Figure 1.17 Fields for a TM<sub>010</sub> mode of a cylindrical (pillbox)-cavity resonator.

## Analytical field solutions: the Pill-box cavity

In the simplest cases the mode field configuration can be calculated analytically, while in almost all practical cases the solutions are computed numerically by means of dedicated computer codes.

One of the most interesting didactical case is the cylindrical or -pill-box" cavity. The pill-box cavity can be seen as a piece of circular waveguide short-circuited at both ends by metallic plates.



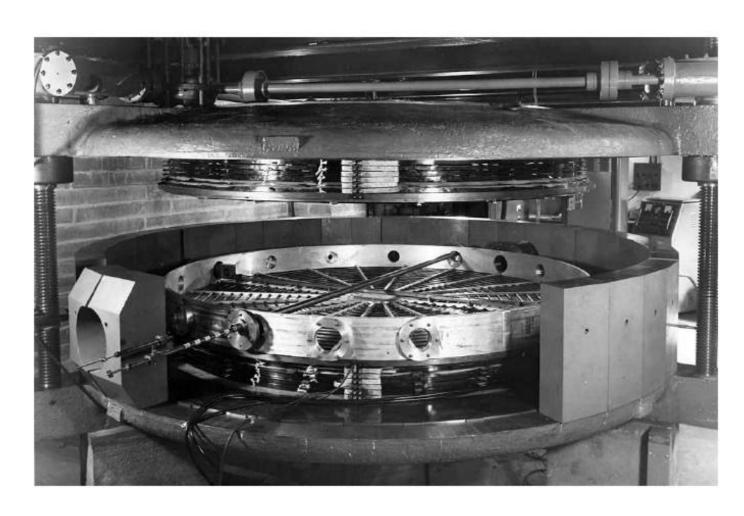
Circular waveguide modes

Wave Type	$TM_{01}$	$TM_{02}$	<i>TM</i> <sub>11</sub>	$TE_{01}$	<i>TE</i> <sub>11</sub>
Field distributions in cross-sectional plane, at plane of maximum transverse fields	ALC: SALE		Distributions Below Along This Plane		Distributions Below Along This Plans
Field distributions along guide					

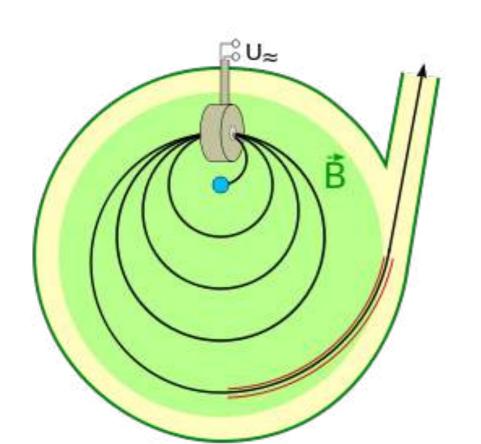


# University College London 28 MeV Microtron at HEP Laboratory





## Microtron - Synchronization



$$t_i = \frac{2\pi R_i}{v_i}$$

$$ev_i B = m_i \frac{v_i^2}{R_i}$$

$$R_i = \frac{v_i m_i c^2}{e c^2 B} = \frac{v_i}{e c^2 B} E_i$$

$$\Delta t = t_{i+1} - t_i = \frac{2\pi}{ec^2B} (E_{i+1} - E_i) = \frac{2\pi}{ec^2B} \Delta E$$

# Energy gain/revolution

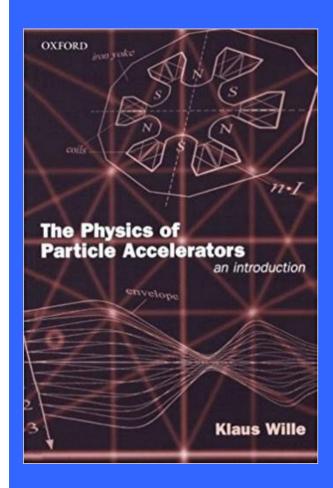
$$\Delta t = t_{i+1} - t_i = \frac{2\pi}{ec^2B} (E_{i+1} - E_i) = \frac{2\pi}{ec^2B} \Delta E$$

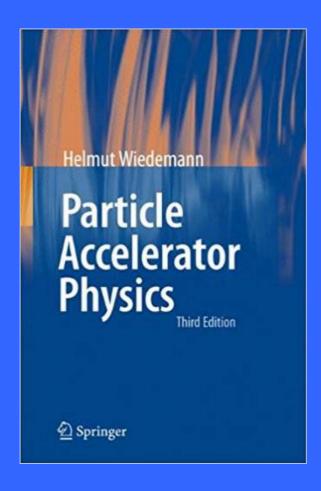
$$\Delta t = \frac{k}{v_{RF}} = \frac{2\pi}{ec^2 B} \Delta E \qquad \Delta E = k \frac{ec^2 B}{2\pi v_{RF}}$$

 In a microtron, due to the electrons' increasing momentum, the particle paths are different for each pass. The time needed for that must be an integer multiple k of the RF period. The allowed energy gain/pass must fulfill the above condition.

## https://www.asimmetrie.it/







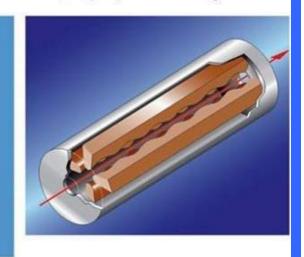
PHYSICS TEXTBOOK

Thomas P. Wangler

**WILEY-VCH** 

### RF Linear Accelerators

Second, Completely Revised and Enlarged Edition





## http://cas.web.cern.ch/previous-schools

2015	Intensity Limitations	Geneva	Switzerland	CERN-2017-006-SP
2015	Advanced Accelerator Physics	Warsaw	Poland	
2015	Accelerators for Medical Applications	Vösendorf	Austria	CERN-2017-004-SP
2014	Plasma Wake Acceleration	Geneva	Switzerland	CERN-2-16-001
2014	Introduction to Accelerator Physics	Prague	Czech Republic	
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2014	Basics of Accelerator Science and Technology at CERN	Chavannes de Bogis	Switzerland	
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2013	Advanced Accelerator Physics	Trondheim	Norway	CERN-2014-009
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