

DEFINITION OF THE LIOUVILLIAN ①

Consider $A(p, q)$ which depends on p, q only
(no explicit time dependence)

Let $p(t), q(t)$ be the solution of the Hamilton equation and consider $A(p(t), q(t))$

$$\begin{aligned} \frac{d}{dt} A(p(t), q(t)) &= \frac{\partial A}{\partial p} \dot{p} + \frac{\partial A}{\partial q} \dot{q} \\ &= \left(-\frac{\partial H}{\partial q} \frac{\partial}{\partial p} + \frac{\partial H}{\partial p} \frac{\partial}{\partial q} \right) A(p(t), q(t)) \end{aligned}$$

The Liouvillian operator is defined as

$$iL = -\frac{\partial H}{\partial q} \frac{\partial}{\partial p} + \frac{\partial H}{\partial p} \frac{\partial}{\partial q}$$

↑
introduced to make L hermitian

$$\frac{d}{dt} A(p(t), q(t)) = iLA$$

If $H(p, q)$ does not depend explicitly on time t .

$$A(p(t), q(t)) = e^{iLt} A(p(0), q(0))$$

↑
 L hermitian $\Rightarrow e^{iLt}$ unitary

[Newton's dynamics is a unitary evolution in phase space

We define the Liouillian op. as

$$\mathcal{L} = \frac{\partial H}{\partial \bar{p}} \cdot \frac{\partial}{\partial q} - \frac{\partial H}{\partial q} \cdot \frac{\partial}{\partial \bar{p}} \quad \left(\begin{array}{l} \text{For 1 particle, for } N \\ \text{particles, add a sum over} \\ \text{particles} \end{array} \right)$$

We wish to show that it is hermitean in the Hilbert space of complex functions of (p, q) with scalar product

$$\langle \psi_1 | \psi_2 \rangle = \int dp dq \psi_1^*(p, q) \psi_2(p, q) \quad \left[\begin{array}{l} \text{We use } \mathcal{M} \\ \text{notation} \end{array} \right]$$

Indeed, consider a 1D system of one particle

$$\langle \psi_1 | \mathcal{L} \psi_2 \rangle = \int dp dq \psi_1^*(p, q) \left(\frac{\partial H}{\partial \bar{p}} \frac{\partial \psi_2}{\partial q} - \frac{\partial H}{\partial q} \frac{\partial \psi_2}{\partial \bar{p}} \right)$$

We integrate by parts and use the fact that $\psi_{1,2}(\pm\infty, q) = 0$ $\psi_{1,2}(p, \pm\infty) = 0$ to guarantee integrability

$$\begin{aligned} &= - \int dp dq \psi_2 \left[\frac{\partial}{\partial q} \left(\frac{\partial H}{\partial \bar{p}} \psi_1^* \right) - \frac{\partial}{\partial \bar{p}} \left(\frac{\partial H}{\partial q} \psi_1^* \right) \right] \\ &= - \int dp dq \psi_2 \left[\cancel{\frac{\partial^2 H}{\partial p \partial q}} \psi_1^* + \frac{\partial H}{\partial \bar{p}} \frac{\partial \psi_1^*}{\partial q} - \cancel{\frac{\partial^2 H}{\partial p \partial q}} \psi_1^* - \frac{\partial H}{\partial q} \frac{\partial \psi_1^*}{\partial \bar{p}} \right] \\ &= - \left[\int dp dq \psi_2^* \left(\frac{\partial H}{\partial \bar{p}} \frac{\partial}{\partial q} - \frac{\partial H}{\partial q} \frac{\partial}{\partial \bar{p}} \right) \psi_1 \right]^* \\ &= - \langle \psi_2 | \mathcal{L} \psi_1 \rangle^* \end{aligned}$$

$$\Rightarrow i \langle \psi_1 | \mathcal{L} \psi_2 \rangle = - (i)^* \langle \psi_2 | \mathcal{L} \psi_1 \rangle^* =$$

$$\langle \psi_1 | \mathcal{L} \psi_2 \rangle = \langle \psi_2 | \mathcal{L} \psi_1 \rangle^* \quad | \quad \mathcal{L} \text{ is hermitian}$$